L 3460-66 ENT(m)/EPA(w)-2/ENA(m)-2 LJP(c) 144 UR/0089/65/018/006/0633/0634 ACCESSION NR: AP5016934 621.384.612 AUTHORS: Vorob'yev, A. A.; Didenko, A. N.; Lisitsyn, Morozov, B. N.; Potekhin, Yu. I.; Salivon, L. G.; TITLE: 10 MeV waveguide synchrotron SOURCE: Atomnaya energiya, v. 18, no. 6, 1965, 633-634 TOPIC TAGS: synchrotron, circular accelerator, electron accelerator, high energy accelerator, waveguide ABSTRACT: After first listing some of the theoretical problems involved in the design of accelerators of this type and dealt with at Institut yadernoy fiziki Tomskogo politekhnicheskogo instituta (Scientific Research Institute of Nuclear Physics of the Tomsk Polytechnic Institute), the authors leser be briefly the 10 MeV synchrotron constructed and in operation at this institute since December 1963. The accelerating system is a rectangular waveguide bent in the shape of a ring, loaded with diaphragms on the outer wall. A standing H<sub>018</sub> mode

L 3460-66

ACCESSION NR: AP5016934

In the π/2 mcde is excited in the waveguide. The radius of the equilibrium orbit of the electrons, on which the phase velocity of the H<sub>018</sub>

wave is equal to the velocity of light, is 13 cm. The waveguide interaction space measures 6 x 6 cm. The system Q is approximately 300, the shunt resistance is approximately 0.07 Meg. The electrons are first accelerated to 3 MeV in the betatron mode by a Kerst gun. The high-frequency electromagnetic oscillations are generated by a pulsed 10-cm generator of 5,000 μsec pulses of 400 W each. The operating presure is 2 x 10<sup>-5</sup> mm Hg. Several of the control and construction features are briefly described. We thank the students of the Tomsk Polytechnic Institute V. I. Zhuravlev, A. M. Voloshin, P. I. Matyazh,

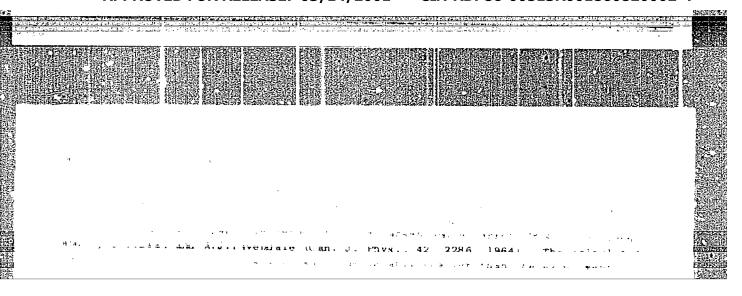
A. A. Kushch, and A. N. Pershin, who participated in the adjustment and startup of the installation, and also Ye. S. Kovalenko and A. P. Ol'shanskiy for participating in the development of the accelerator theory. Its design, and model test.' Orig. art. has: I figure

ASSOCIATION: None

Card 2/3

L 3460-66
ACCESSION NR: AP5016934

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Eighty-meter sectional belt conveyer. Mekh.stroi. 4 no.10:8-11
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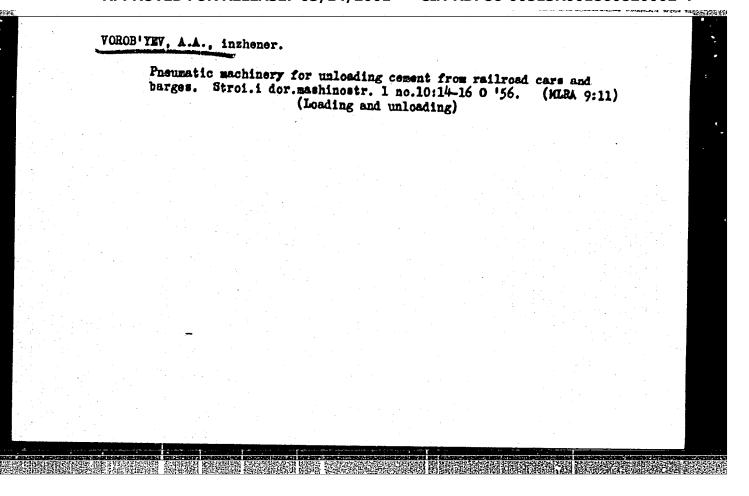
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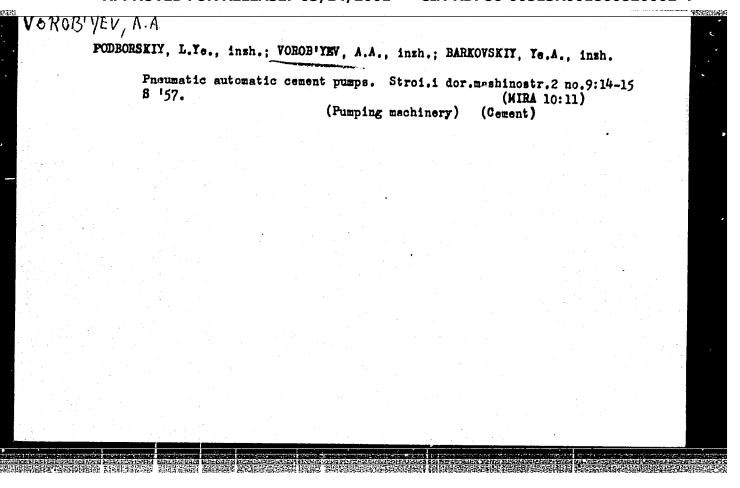
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RASTVOROMASOS PROIZVODITEL'NOST'YU 1-2 m3/chas DLYA ZHESTKIKH RASTVOROV Page 143

SO: Collections of Annotations of Scientific Research Work on Construction, completed in 1950. Moscow





VOROB! YEV, A., inzh.								
Production of "silicalcite" products is being organized in Moldavia. Stroi.mat. 3 no.11:40 N '57. (MIRA 10:12 (Kishinev-Silicates-Congresses)							10:12)	
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VOROB'YEY, A., inzh.

Mechanized assembly line production of shell rock, Stroi. mat. 4
no.9:12-13 S '58. (MIRA 11:12)

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VYGODCHIKOV, G.V.; GEKKER, V.D.; LARINA, I.A.; SERGEYEVA, N.S.; VOROB'YEV, A.A.; SALTIKOV, R.A.

Hasic principles underlying the production of polyvalent vaccines against anaerobic and intestinal infections. Zhur. mikrobiol., epid. i immun. 40 no.3:9-14 Mr \*63. (MIRA 17:2)

1. Iz Instituta epidemiologii i mikrobiologii imeni Gamalei AMN SSSR.

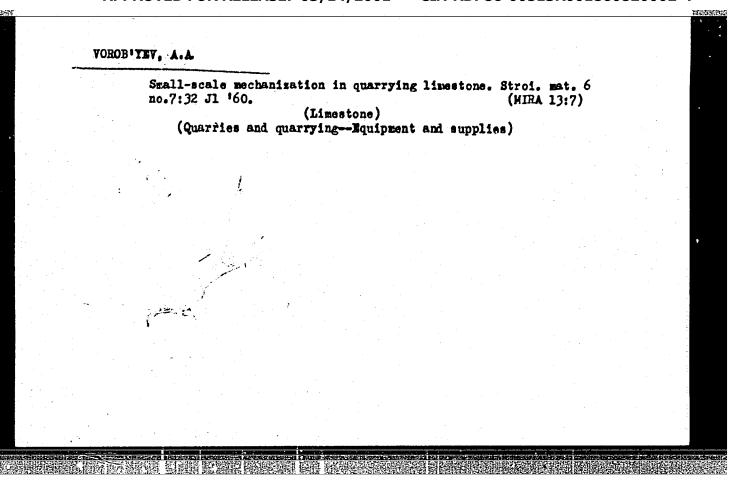
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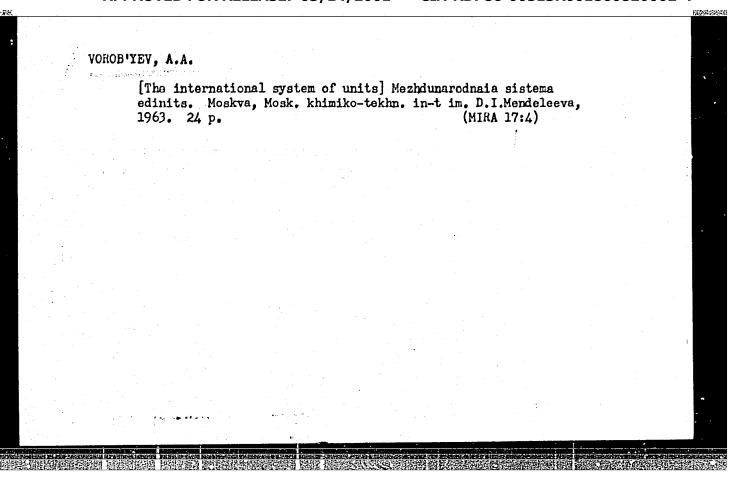
Using granulated slags in making high-strength concretes for prestressed fleer panels. Bet. i zhel.-bet. no.6:234-235 Je '58.

(Kishinev--Concrete)

# VOROB'YEV, A, Over-all mechanization of cement unloading from railroad cars. Stroitel' no.11:23 '58. (MIRA 11:12) 1. Rukovoditel' laboratorii oborudovaniya dlya razgruzki i pnevmaticheskogo transporta tsementa Leningradskogo filiala Vsesoyuznogo nauchnoisəledovatel'skogo instituta stroitel'nogoi dorozhnogo mashinostroyeniya. (Loading and unloading) (Cement--Transportation)

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VOROBYEV, 11.H

AUTHORS:

Bochagov, B. A., Vorob'yev, A. A., Komar, A. P. 57-27-7-20/40

TITLE:

An Impulse Ionization Chamber as a Device for the Simultaneous Investigation of the Energetic and Angular Distributions of Charged Particles (Impul'snaya innizatsionnaya kamera kak pribor alya odnovremennogo izucheniya energeticheskikh i uglovykh raspredeleniy zaryazhennykh chastits).

PERIODICAL:

Zhurnal Tekhnicheskoy Fiziki, 1957, Vol. 27, Nr 7, pp. 1575-1577

(ussa)

ABSTRACT:

It is shown that the energy E (half width of the lines of a-spectra) and the angle of flight  $\phi$  (between the normal to the electrodes and the direction of flight of the charged particle) of the particle concerned can be determined beginning from the source, when the quantity of the impulse  $V_1$  (the voltage at the collecting electrode) and the quantity of one of theimpulses  $V_2$  (the voltage at the high-voltage electrode),  $V_3$  (the voltage at the power supply or  $V_4$  (the voltage at the moment where all electrons have reached the collecting electrode) is simultaneously measured. The accuracy of measurement of  $\cos \phi$  in this connection is about 3% and can be improved. At present a mechanical collimator is often used in measurements of the energy of a-particles may also be brought to collimation without a mechanical collimator due to the fact that the ionization chamber permits a minultaneous measurement of E and

Card 1/2

APPROVED FOR RELEASE: 03/14/2001 CIA-RDP86-00513R001860820002-4"

An Impulse Ionization Chamber as a Device for the Simultaneous 57-27-7-20/40 Investigation of the Energetic and Angular Distributions of Charged Particles.

> φ. The method suggested here can also be employed for the solution of problems of a-spectroscopy, for theinvestigation of the α- - correlation, the neutron-spectrum according to the protons given off and for the investigation of the angular distribution of heavy products of nuclear reactions. There are 3 figures.

ASSOCIATION:

Leningrad Polytechnic Institute imeni M. I. Kalinin (Leningrads-

kiy politekhnicheskiy institut im. M. I. Kalinina)

SUBMITTED:

January 27, 1956

AVAILABLE:

Library of Congress

1. Ionization chambers-Applications 2. Particles-Energy-Measurement

3. Particles-Transmission-Analysis

Card 2/2

### "APPROVED FOR RELEASE: 03/14/2001

CIA-RDP86-00513R001860820002-4

PA - 2994 YOROB'YEW, A.A., KOROLEV, V.A., KOMAR, A.P., AUTHOR

SELIVERSTOV, D.E., The Coefficient of the Interior Conversion of  $\gamma$ -Radiation with the · TITLE

Energy 53 KeV on the L-shell of the Th 430.

(Koeffitsyent vnutrenney konversii γ-izlucheniya energii 53 keV na

L-oboloche Thaso - Russian)

Zhurnal Eksperim. i Teoret. Fiziki, 1957, Vol 32, Nr 3, PERIODICAL

pp 623-623, (U.S.S.R.) Reviewed 7/1957 Received 6/1957

According, to the data obtained from publications this coefficient is pro-ABSTRACT bably large. Theauthors determined this conversion coefficient by means of the method of a-y coincidences. An enriched UR34 source was used. The aparticles were recorded by means of a momentum ionization chamber and the Y-quanta by means of a scintillation counter with an NaJ(T1)-crystal. The  $\gamma$ -spectrum was recorded in coincidence with the  $\alpha$ -particles which lead to the basic level and to the first excited level of the Th230. This radiation originates entirely from the inner conversion on the L-shells of the Th<sup>230</sup>. The coefficient of conversion was determined from the ratio of the number by of the radio X-ray quanta (without absorber) to the number Ny of 53 keV quanta (which were reduced to the same number No of the recorded a-particles.) The result Nn/N. = 13c was obtained. The error committed in measuring: remains below 50%. The extrapolation of the theoretical value furnishes the following values for the sum of the coefficients of conversion on the

LI-,LII- and LIII shells, according to the type of radiation, Card 1/2

### "APPROVED FOR RELEASE: 03/14/2001

CIA-RDP86-00513R001860820002-4

The Coefficient of the Interior Conversion of  $\gamma$ -Radiation with the Energy 53 KeV on the L-Shell of the Th<sup>230</sup>.

 $E_1$   $E_2$   $E_3$   $M_1$   $M_2$  (1,0) 170 (1,0) (1,0) (1,0) (1,0)

A comparison with experimental results permits the conclusion that the radiation observed is an electric quadrupole radiation. Because the ground state of the even-even nuclei has the angular momentum 0 and the parity +, the first excited level of Th<sup>230</sup> must have the angular momentum 2 and the parity +. The results obtained experimentally confirm the rotation-like nature of this level (corresponding to BOHR'S model).

ASSOCIATION PRESENTED BY SUBMITTED AVAILABLE

Card 2/2

Leningrad Physical-Technical Institute of the Academy of Science of the

17.12.1956. Library of Congress.

SOV/120-59-1-21/50

AUTHORS: Vorob'yev, A. A., Korolev, V. A., Solyakin, G. Ye.

TITLE: Measurement of the Grid Current in the Tubes Employed in Low-Noise Amplifiers (Izmereniye setochnogo toka v lampakh, ispol'zuyemykh v usilitelyakh s nizkim shumom)

PERIODICAL: Pribory i tekhniku eksperimenta, 1959, Nr 1, pp 85-89 (USSR)

ABSTRACT: It is known from the Nyquist theory that the noise produced by the grid current can be expressed by:

$$\frac{\overline{U_{\text{sh.s.}}^2} = \overline{u_c} \int_0^2 \frac{e^2 F(\omega) d\omega}{1 + \omega^2 \tau^2}, \qquad (3)$$

where  $\tau$  = RC, I<sub>C</sub> is the grid current; R is the grid leak of the tube and C is its input capacitance; function  $F(\omega)$  in Eq (3) is formed by the product of the transfer functions of an integrating and a differentiating network; the time constants of the networks are  $T_1 = T_2 = T$ . Consequently,

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# SOV/120-59-1-21/50

Measurement of the Grid Current in the Tubes Employed in Low-Noise Amplifiers

the grid current noise can be expressed by Eq\_(4), where q = 7/T . If the tube contains a resistor R at its input, the noise due to this can be expressed by Eq (6). Provided the noise due to this can be expressed by Eq. (6). Frowtight the same function  $F(\omega)$  is used, the integration of Eq. (6) results in Eq. (7). The maximum value of the thermal noise, expressed by Eq. (7), occurs when the resistance is given by Eq. (8); this value is given by Eq. (9). On the other hand, the maximum value of the noise produced by the grid current is given by Eq (5). Consequently, the grid current can be expressed in terms of a ratio of the maximum grid current noise to the thermal noise and this is expressed by Eq (10). This equation can be used for determining the value of  $I_{\mathbb{C}}$ . By comparing Eqs (3) and (6), it is found that the relationship between the grid current noise and the thermal noise is expressed by Eq (15). This can also be used for determining ; for example, if a value of R is determined such that the current noise is equal to the thermal noise, the grid is the value of R current is given by Eq (16); here, Ro necessary to secure the equality of the two noises. The above

Card 2/3 methods were employed to measure the grid current in the tube

SOV/120-59-1-21/50

Measurement of the Grid Current in the Tubes Employed in Low-Noise Amplifiers

Type 6ZhlP which were operated as triodes with an anode voltage of 60 V and a heater voltage of 6 V. The dependence of the total noise on the input resistance is illustrated in Fig 2. From this it is found that the grid current was 1.0. 10<sup>-10</sup>A, when determined from Eq (11) (or from Eq 14), and it was 1.15 x 10<sup>-10</sup>A when evaluated from Eq (16). The authors express their gratitude to F. M. Sobolevskaya for her help in the measurements, to S. N. Nikolayev for discussing the results, and to A. P. Komar for his interest in this work. The paper contains 3 figures and 2 references, of which 1 is English and 1 is Soviet.

ASSOCIATION: Leningradskiy fiziko-tekhnicheskiy institut AN SSSR (Leningrad Physics Engineering Institute of the Soviet Academy of Sciences)

SUBMITTED: February 5, 1958.

Card 3/3

SOV/120-59-2-27/50

AUTHORS: Vorob'yev, A.A., Korolev, V.A. and Solyakin, G.Ye.

TITLE: The Choice of Optimum Pass-band in an Amplifier Working with an Ionization Chamber (Vybor optimal noy polosy

propuskaniya v usilitele, rabotayushchem s

ionizatsionnoy kameroy)

PERIODICAL: Pribory i tekhnika eksperimenta, 1959, Nr 2, pp 95-102 (USSR)

ABSTRACT: A calculation is made of the optimum bandwidth of an amplifier with two differentiating circuits. It is shown that the introduction of the second differentiating circuit completely avoids the influence of microphonic

effects and low frequency noise without deteriorating the signal-to-noise ratio. The resolving power of an ionization alpha-spectrometer is determined basically by

the noise in the first valve. When the leakage resistance of the first valve is high enough thermal noise may be neglected and only the contributions of anode and grid current taken into account. Usually the

anode and grid current taken into account. Usually the maximum signal-to-noise ratio is guaranteed by correct choice of amplifier bandwidth and this usually means the state of the state o

Card 1/4 inserting a differentiating and an integrating circuit.
This case has already been considered by Elmore in Ref 1.

SOV/120-59-2-27/50 The Choice of Optimum Pass-band in an Amplifier Working with an Ionization Chamber

This scheme has a number of drawbacks; in particular the location of the differentiating circuit is difficult, since it is preferable to place it before the amplifier in order to avoid overloading on microphony, but also convenient to place the circuit within the middle of the amplifier when A.C. heaters are used. In the analysis for brevity an arrangement of one differentiator followed by one integrator is described as {1,1}; the cases {1,2} {2,2} are also considered. The spectral densities of the grid and anode currents are given by Eqs (1) and (2). For the three circuit combinations described above, expressions for the minimum value of noise are given by Eqs (8), (12) and (17). In the many curves which are presented two parameters are used; p which is the ratio of the time constants of the integrator and the differentiator circuits, and a which is defined in In calculating signal-to-noise ratio it is assumed that a rectangular voltage pulse is delivered Card 2/4 from the ionization chamber. Signal-to-noise ratio is denoted by Q. In Fig 1 the signal-to-noise ratio is

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SOV/120-59-2-27/50

The Choice of Optimum Pass-band in an Amplifier Working with an Ionization Chamber

given by a solid line and the signal amplitude by the dotted lines. Figs 3, 4 and 5 show for the three circuit arrangements respectively the variation of signal-to-noise ratio with p for various pulse durations. Figs 6, 7 and 8 are the corresponding figures with p and a as parameters. Ionization Ionization chambers suffer from microphony at frequencies up to 100 c/s. By using two differentiating circuits the contribution to the microphony may be reduced with respect to that due to valve noise by a factor of approximately 100 at a frequency of 100 c/s; at lower frequencies this reduction is even more significant. It has so far been assumed that the voltage pulses are truly rectangular; in practice they have sloping fronts and if these slopes are linear it is possible to calculate easily the loss in amplitude as a function of the differentiating and integrating circuits. This loss is shown plotted in Figs 9 and 10 respectively for single and double circuits. Table 1 summarizes the amplitude loss for various rise times for the three types

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SOV/120-59-2-27/50 The Choice of Optimum Pass-band in an Amplifier Working with an Ionization Chamber

of circuit; this is experimental data. For all three circuits the relationship between amplitude loss and rise time is quadratic. In Table 2 experimental and calculated results are compared for various values of differentiator and integrator time constant; this table applies to the case of {1,2}. The authors thank

Card 4/4 M.F. Sobolevakaya and A.P. Komar.

There are 10 figures, 2 tables and 2 English references.

ASSOCIATION: Fiziko-tekhnicheskiy institut AN SSSR (Physico-Technical Institute of the Academy of Sciences,

USSR)

SUBMITTED: February 13, 1958

### "APPROVED FOR RELEASE: 03/14/2001

### CIA-RDP86-00513R001860820002-4

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Vorob'yev, A. A., Ivanov, B. A., Komar, A. P., Korolev, V. A.

TITLE:

**AUTHORS:** 

Influence of Ramsnuer-Townsend Effect on the Mobility of Electrons

in Spectroscopically Pure Argon

PERIODICAL:

Zhurnal tekhnicheskoy fiziki, 1959, Vol 29, Nr 10, pp 1252-1258

(USSR)

ABSTRACT:

The purpose of the paper is to verify the influence of Ramsauer-Townsend effect on the mobility of electrons in spectroscopically pure argon. The study is experimental in nature. The drift of electrons is measured as a function of E/p, where E is intensity of the electric field and p is barometric pressure of argon in the experimental chamber. The experiments were carried out for values of E/p between 0.001 and 1.5. At small values of E/p a maximum was observed similar to that obtained by other investigators. This maximum could be explained as the result of the Ramsauer-Townsend effect, or it might have been caused by the excitation of molecular levels owing to the presence of impurities in argon. For this reason industrial argon of 99.6% was also used. The ionization chamber was filled with argon at

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Influence of Ramsauer-Townsend Effect on the Mobility of Electrons in Spectroscopically Pure Argon

75333 sov/57-29-10-10/18

pressures of up to 1,000 mm Hg. For a source of (&-particles U<sup>234</sup> was used. The impurities in the spectroscopically pure argon amounted to less than 10<sup>-14</sup>%. The total sum of errors did not exceed 20%. The mean free path & of the electrons at a pressure of 1 mm Hg and the average fraction of energy f lost in one collision were measured, and the values obtained fully corroborated the influence of the Ramsauer-Townsend effect on the election mobility in the spectroscopically pure argon. When industrial argon was used it was found that the value of f is affected by argon impurities influencing the maximum value of E/p. This, however, is not to be taken as the result of Ramsauer-Townsend effect. There are o figures, and 3 references, 6 U.S., 1 Canadian, and 1 French. The U.S. and Canadian references are: Nielsen, R. A., Phys. Rev.; 50, 950, 1936; Klema, E. D., Allen, J. S., Phys. Rev., 77, 661, 1950; Kirshner, J. M., Toffolo, D. S., J. Appl. Phys., 23, 594, 1952; Bortner, T. E., Hurst, G. S., Stone, W. G., RS1, 28, 103, 1957; Bell, P.R., Jordan, W. H., Davis, R. C., Phys. Rev., 83, 490, 1951; Colli, L., Facchini, U., RS1, 23, 39, 1952; English, W. N., Hanna, G. C., Can. J. Phys., 31, 763, 1953.

Card 2/3

Influence of Ramsauer-Townsend Effect on the Mobility

of Electrons in Spectroscopically Pure Argon

793 ;3 \$69/57-70-10-10/18

ASSOCIATION:

Institute for Technical Physics, Academy of Sciences, USSR

(Fiziko-tekhnicheskiy institut, AN SSSR)

SUBMITTED;

April 27, 1959

Card 3/3

24(7) AUTHORS:

Vorob'yev, A. A., Komar, A. P., Korolev, V. A., SOV/56-37-2-32/56

Solyakin, G. Ye.

TITLE:

The &-Spectrum of the Natural Mixture of Isotopic Samarium

PERIODICAL:

Zhurnal eksperimental noy i teoreticheskoy fiziki, 1959,

Vol 37, Nr 2(8), pp 546 - 548 (USSR)

ABSTRACT:

In the present "Letter to the Editor" the authors report on investigations of the a-spectrum of Sm 147 and the isotopic mixture by means of a pulse ionization chamber; the chamber was filled with chemically pure argon (99.9% Ar, +02, N2,002). The measured a-spectrum of Sm 147 is shown by figure 1; it has a half width of 43 kev (when intensive a-emitters, as e.g., U234, were used, the half width amounted to 30 keV). The energy

of the  $\alpha$ -particles of Sm 147 was determined as amounting to (2.19 ± 0.01) MeV, which agrees well with the value mentioned in reference 6. Figure 2 shows the spectrum of the  $\alpha$ -particles of the natural isotopic mixture (without collimation) within

Card 1/2

The a-Spectrum of the Natural Mixture of Isotopic Sumarium SOV/56-37-2-32/56

the energy interval of 2.0 - 2.8 Mev. The energy of the co-particles of  $\rm Sm^{146}$  is (according to reference 7) equal to  $\sim 2.55$  Nev; knowledge of this fact and of the entire background (within the range of 1.5 - 2.5 Mev - 1 pulse/hour) makes it possible to evaluate the upper limit of the Sm 146 content in the natural isotopic mixture and thus to determine the half lives:  $T(Sm^{147}) = 10^{12}a$  and  $T(Sm^{146})$  -= 5.107a. The Sm 146-concentration in the natural isotopic mixture is not greater than 2.5.10-6 % (the number of d-particles originating from Sm 146-decay does not exceed the background). According to a mass-spectrometric analysis the content would amount to 8.10-5 % (Ref 8). There are 2 figures and 8 references, 1 of which is Soviet. Leningradskiy fiziko-tekhnicheskiy institut Akademii nauk SSSR

ASSOCIATION:

(Leningrad Physico-technical Institute of the Academy of

Sciences, USSR)

SUBMITTED: Card 2/2

March 26, 1959

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Fig. 2 Book Exponential 607/533)  Fig. 2 Book Exponential 607/533)  Fig. 2 Book Exponential 607/533)  Fig. 2 Book Exponential 607/533  Fig. 6 Fig. 7 Fig. 6 Fig. 6 Fig. 6 Fig. 7 Fig	Anniar Sin, Richard 12, Daitrive, Ja. A. Teplove, sed L.F. Fatrove.	GENERATION OF MAINTINES TO ANNI TO STATE OF THE PROPERTY OF TH	Ecty V.I., A.B. Eggetsow, and R.D. Rulte. Effect of Baltiche Eighteife and Radation Daring Riverson Pailsbup in Acceleration			
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s/089/60/009/004/010/020 В006/В070 Voroblyev, A. A., Pokrovskiy, S. F. Comparison of Circuits of Cascade Generators for the Production of Large Currents With Small Pulsation AUTHORS: Atomnaya energiya, 1960, Vol. 9, No. 4, pp. 305 .. 308 TITLE: TEXT: For making exact measurements on thin targets in the energy range of 2 - 3 MeV, it is necessary that the particle energy be constant up to 0.1 - 0.5% for a beam current of the order of 5 - 10 ma. The authors have PERIODICAL: investigated the possibility of using for this purpose an electrostatic accelerator which works with a cascade generator of a power of up to 30 kw. They have compared the calculations of cascade-generator circuits according to different theories with one another and with experimental data. The formulas of calculations for four cascade generators are collected in Table 1. They are taken from papers of V.S. Melikhov. V. S. Novikovskiy, A. Bouwers, and the authors of the present "Letter to the Editor". Fig. 1 shows the dependence of the pulsation of the output notential All on the number of stages n for a cascade generator with a potential by on the number of stages n for a cascade generator with a Card 1/3

81,230

8/089/60/009/004/010/020 B006/B070

Comparison of Circuits of Cascade Generators for the Production of Large Currents With Small Pulsation

selenium rectifier, which has a charging current of 1 ma, an input potential of frequency 50 cps, and 2 - 7 stages. Fig. 2 shows the dependence of the fall in the output potential on n for a charging current of 2 ma for one and the same generator circuit. The results of comparison between the theoretical and experimental results are summarized as follows (the methods of calculation are named after their authors): 1) For a Cockcroft-Walton generator, the best agreement is obtained by using the formulas of Vorob'yev and Melikhov; 2) for calculating the voltage pulsation for a symmetric circuit, the best formula is that of Novikovskiy, and for a more exact calculation, that of Pokrovskiy; for calculating the fall of potential, the best formula is that of Novikovskiy and Pokrovskiy; 3) for a three-phase circuit, the best formula is that of Pokrovskiy. The properties of four types of cascade generator are compared in Table 2. This leads to the following conclusions; 1) The Cockcroft-Walton circuit is the simplest, but pulsation and voltage fall are large. 2) Cascade-generator circuits with capacitances which decrease linearly with the distance from the source of potential have 13 times less pulsation and 17 times less fall of

Card 2/3

Comparison of Circuits of Cascade Generators for the Production of Large Currents With Small Pulsation S/089/60/009/004/010/020 B006/B070

potential compared to the Cockcroft-Walton circuit. 3) Symmetric circuits have 26 times less pulsation and a fourth of the potential fall compared to the Cockcroft-Walton circuit, and so these circuits have more rectifiers and capacitors. 4) The three-phase circuit is more complicated than the symmetric one, and has a smaller potential stability, more capacitors and rectifiers, but has the highest charging current possible (three times larger than that for the Cockcroft-Walton generator). There are 2 figures, 2 tables, and 3 references: 2 Soviet and 1 German.



SUBMITTED: March 4, 1960

Card 3/3

55/1/1/1

B/089/60/009/006/004/011 B102/B212

21,3000

AUTHORS:

Surkov, Yu. A., Vorob'yev, A. A., Korolev, V. A.,

Vilenskiy, V. D.

TITLE: Investigation of the composition of urarium isotopes in rare-

earth minerals

PERIODICAL: Atomnaya energiya, v. 9, no. 6, 1960, 477-482

TEXT: The authors have tried to find out whether the isotope  $Cm^{247}$  exists (or existed) in nature (it is produced during plutonium irradiation in a reactor). This isotope changes over into  $Pu^{243}$  with a half-life of the corder of  $10^8$  years and finally into  $U^{235}$ . One may assume the following reaction chain  $Cm^{247} \xrightarrow{\alpha} Pu^{243} \xrightarrow{\beta} Am^{243} \xrightarrow{\alpha} Pu^{239} \xrightarrow{\beta} Pu^{239} Pu$ 

 $\alpha$   $U^{235}$  ..., from the ratio  $U^{235}/U^{238}$  one could conclude that there still exists  $Cm^{247}$  in very old rare-earth minerals. The authors investigated the composition of uranium isotopes in xenotime, orthite and

Card 1/8

S/089/60/009/006/004/011 B102/B212

Investigation of the ...

gadolinite with an age of 2.109 years. The samples had been furnished from the Mineralogicheskiy muzey AN SSSR (Mineralogy Museum of the AS USSR). The uranium was separated radiochemically from the minerals for an  $\alpha$ -spectrometric analysis. The relative content of U<sup>238</sup> and U<sup>238</sup> was determined from the  $\alpha$ -activity of these isotopes. An ionization chamber with screen (see Fig.2) had to be utilized since the uranium content was minute (0.25 - 1 mg). The chamber was filled with Ar + 0.5% CH<sub>4</sub>; the  $\alpha$ -radiating preparation was located on the high-voltage electrode. The  $\alpha$ -particles will hit the collector electrode with a time delay of t<sub>delay</sub>=(d-Rcosp)/w according to their direction of flight; R denotes the range of the  $\alpha$ -particles, w the electron drift rate, d the distance between high-voltage electrode and acreen,  $\phi$  the angle between the direction of flight of the  $\alpha$ -particle and the normal. The nethod of time collimation applied for the purpose consists in that only those pulses are recorded, for which tdelay < t'; thus, the pulses from  $\alpha$ -particles emitted at small angles were eliminated. The degree of collimation was characterized by f (f/w =  $t_{max}$ -t¹). The share  $\varrho$  of the recorded pulses from a-particles is given by  $Q = 1-f/R = N/N_0$ , where N<sub>0</sub> Card 2/8

# Investigation of the... S/089/60/009/006/004/011 B102/B212Investigation of the... and N denote the intensities of a line before and after the collimation, and N denote the intensities of a line before and after the collimation, and N denote the intensities $N_1^0$ and $N_2^0$ ) and $Q^0 = \frac{N_1^1}{N_1^2} = \frac{N_1^1}{N_2} = \frac{1}{N_1^2} = \frac{Q_1^1}{1 - \frac{1}{R_1}} = QA, \quad (4).$ Practically, there were three lines for the unanium isotopes. $Q^0_{13} = \frac{N_{131}^1}{N_{131}^2}, \quad Q_{234} = \frac{N_{132}^1}{N_{131}^2}, \quad Q_{235}^1 = \frac{N_{132}^1}{N_{132}^2}.$ $Q_{235} = \frac{N_{131}^1}{N_{131}^2}$ The following holds $Q^0_{235} = Q_{233}A_{233} = Q_{235}Q_{235} = \frac{Q_{235}}{Q_{235}} = \frac{Q_{235}}{N_{132}^2}.$ $= Q_{233}\left(1 + \frac{Q_{137} - Q_{133}}{Q_{235}}\right) = Q_{235}\left(1 + P_{235}\right), \quad (5)$ and analogously $Q^0_{234} = Q_{234}A_{234} = Q_{234}\left(1 + P_{234}\right), \quad (6)$ $P_{234} = \frac{1}{1 - \frac{1}{R_{131}}}\left(\frac{1}{R_{131}} - \frac{1}{R_{131}}\right). \quad (6a)$ Card $3/\frac{8}{3}$

Investigation of the...  $P_{235} \text{ and } P_{234} \text{ are interrelated by}$   $P_{235} = \frac{1}{R_{13}} \frac{1}{R_{13}} \frac{1}{R_{13}} P_{234} = \frac{1}{R_{23}} \frac{1}{R_{23}} \frac{1}{R_{23}} \frac{1}{R_{23}} P_{234} = \frac{1}{R_{13}} \frac{1}{R_{13}} \frac{1}{R_{13}} \frac{1}{R_{13}} P_{234} = \frac{1}{R_{23}} \frac{1}$ 

Investigation of the..

5/089/60/009/006/004/011 B102/B212

small. It was possible to determine  $U_{235}/U_{238}$  with an error of ~2%. The test data are compiled in Table 1 (without collimation) and Table 2 (with collimation). It is apparent that the ratio of the isotopes was a little higher in gadolinite ( $q \approx 1.046 \pm 0.02$ ). Here, it may be assumed that this raise is due to the existence of  $Cm^{247}$ . If its half-life is taken as  $\approx 4.10^7$ a then it is possible to calculate the initial content of  $Cm^{247}$  in gadolinite (at a mean uranium content of 0.06%) which could have been  $\approx 10^{-3}$ %. The authors thank A. P. Komar and V. I. Baranov for their interest in these investigations. There are 4 figures, 2 tables, and 7 references: 4 Soviet-bloc and 3 non-Soviet-bloc.

SUBMITTED: February 24, 1960

Card 5/8

**63672** s/048/60/024/009/005/015 B013/B063 Korolev, V. A. The Operation of the Ionization Alpha Spectrometer 21,5300 AUTHORS: Izvestiya Akademii nauk SSSR. Seriya fizicheskaya, 1960, High Loads TITLE: TEXT: In the present paper, the authors suggest a method of raising the Vol. 24, No. 9, pp. 1086 - 1091 permissible loading of ionization chambers up to 10 pulses/sec. The PERIODICAL: permissible loaning of louring in the ionization chamber depends on the duration of the pulses occurring in the ionization of the pulses occurring in the division of the pulses occurring in the ionization occurri duration of the pulses occurring in the following 10 : 15% of methane to duration of electron accumulation in it. By adding 10 : 15% of methane to the engage and by shortening the distance a down to 20.5 cm it is nossible the argon and by shortening the distance q down to ~0.5 cm it is possible the argon and by shortening the distance q down to  $\sim 0.7$  cm  $^{11}$  18 pubblished to shorten the duration of the build-up of pulses down to 0.5-0.4  $\mu sec.$ It is apparently convenient to shape pulses by means of two short-circuited circuits. This method is particularly advantageous because pulse tails can be prevented and the constant component of the pulse spectrum is almost Zero. The authors thoroughly studied the effect of the rise time on the Card 1/2

CIA-RDP86-00513R001860820002-4" APPROVED FOR RELEASE: 03/14/2001

The Operation of the Ionization Alpha Spectrometer Under High Loads

S/048/60/024/009/005/015 B013/B063

pulse height (Fig. 1) and the signal-to-noise ratio (Figs. 2 and 3). The eperation of the spectrometer under high load was checked by measuring the α-spectrum of a Pu<sup>238</sup> source with an intensity of 8 · 10<sup>3</sup> pulses/sec. Besides, the spectrum of generator pulses was measured, which were supplied to the input of the amplifier together with the pulses of α-particles (Figs. 4 and 5). The α-spectrum of Pu<sup>238</sup> is shown in Fig. 6. Its half-width noise makes the largest contribution to the half-width line (45 kev). This noise can be largely reduced by switching on additional tubes. There is reason to believe that the half-width of the α-line can be reduced to at least 40 ÷ 50 kev. Some applications of a strongly loaded spectrometer are for help in the work. There are 6 figures and 3 references: 2 Soviet and

ASSOCIATION: Fiziko-tekhnicheskiy institut Akademii nauk SSSR (Institute of Physics and Technology of the Academy of Sciences USSR)

Card 2/2

5/048/60/024/009/006/015 BO(3/BO53

24.6810 AUTHORS:

Vorch'yev, A. A., Komar, A. P., Korolev, V. A.

TITLE:

Investigation of the Alpha Decay of U235 by Means of an THE PROPERTY OF THE PARTY OF TH Ionization Alpha Spectrometer

Izvestiya Akademii nauk SSSR. Seriya fizicheskaya, 1960,

PERIODICAL:

Vol. 24, No. 9, pp. 1092-1098

The present paper contains the results of an investigation of the d-decay of U<sup>235</sup>, which was carried out by means of an ionization d-spectrometer. The spectrometer was tuned to a 8-spectrometer. The authors studied a spectrum without coincidence with w-quanta (Figs. 1 and 2) and a series of r-spectra coinciding with different groups of alpha particles (Figs. 3 and 4). A source enriched in U235 was used for the measurement. The spectral line of U<sup>234</sup>(98%) showed, however, the highest intensity. It was used to stabilize the amplification factor of the amplifying part of the &-spectrometer. Besides, this line served as a standard for the

Card 1/2

Investigation of the Alpha Decay of U<sup>235</sup> by Means of an Ionization Alpha Spectrometer

S/048/60/024/009/006/015 B003/B063.

measurement of the energies of  $\mathbf{U}^{235}$  alphas. The energies, intensities, and forbiddances of the alpha groups are given in Table 1. The results of the analysis of the U235 d-spectrum agree with the results of Ghiorso although the latter are only of a qualitative character. The a-spectrum of  ${\tt U}^{235}$  has been recently studied by S. A. Baranov and A. G. Zelenkov by means of a spectrometer of high luminous power. The energies of the lines they found are fairly consistent with the data obtained by the present authors. Table 2 gives the results of the determination of multipole or g-transitions. On the basis of measurements of d- and gspectra, the authors suggest a possible d-decay scheme of U235 (Fig. 5). The levels were identified with the help of Nilsson's scheme. Though this identification cannot make a claim to finality, it does not contradict the experimental data available at present. The authors thank S. A. Baranov and A. G. Zelenkov for discussions and information, as well as M. F. Sobolevskaya for her assistance in the measurements. There are 5 figures, 2 tables, and 9 references: 3 Soviet, 3 US, and 1 Danish.

Card 2/2

vorobyEV, A.A. 5/056/60/039/01/11/029 82602 B006/B070 Vorob'yev, A. A., Komar, A. P., Korolev, V. A. Measurement of the Energy of the a-Particles of an Emitter 24.6520 Zhurnal eksperimental'noy i teoreticheskoy fiziki, 1960, AUTHORS: TEXT: The authors measured the α particle energies with the help of an TITLE: Comparing the pulse appetrometer. The ionization was determined by comparing the pulse PERIODICAL: heights of the & particles with those of the generated pulses whose amplitude could be measured to an accuracy of ~ 0.01%. The chamber used ampritude could be measured to an accuracy of the & line was 35 kev.

was filled with 97% Ar and 3% CH<sub>4</sub>. The width of the & 228 Table 1 collects a number of relevant data. The was chosen as a standard. In the first column of this table the energy delices deduced by Table 1 collects a number of relevant data. The was chosen as a stan ard. In the first column of this table the energy values deduced by a magnetic analysis are given the second column of the second magnetic analysis are given, the second column gives ionization I, and magnetic analysis are given, the second column gives lonization 1, and the third the energy calculated according to the formula (1)  $(E_{\rm st}-84)/(E_{\alpha}-84) = I_{\rm st}/I_{\alpha}, \text{ where } E_{\rm st} \text{ and } E_{\alpha} \text{ denote the } \alpha \text{ energies}$ of the standard and the emitter under investigation. There is good Card 1/3

Measurement of the Energy of the a-Particles of an Emitter

\$/056/60/039/01/11/029 B006/B070

agreement between the first and the third columns, from which the conclusion is drawn that the method of energy determination from the ionization in the a spectrometer is suitable. The following results are obtained:

 $At^{217}$ : 7.064 ± 0.005

Fr<sup>221</sup>: 6.336 <u>+</u> 0.005  $v_{\rm T}^{235}$ : 4.396 ± 0.003

Po<sup>213</sup>: 8.368 ± 0.010

 $v^{238}$  . 4.190  $\pm$  0.005

 $v_{II}^{235}$ : 4.211  $\pm$  0.003 These values are compared with the results obtained by other authors. Agreement is good in some cases and not so good in others. Some particular cases in this connection are discussed. Thus, for example, the values

obtained for the two intensive U235 lines (I and II) diverge considerably from those obtained by magnetic spectrometer (Ref. 6). In connection with this, it is pointed out that the measurements lately made by S. A. Baranov, A. G. Zelenkov et al. (Ref. 8) of the  $\alpha$  spectrum of  $U^2$ 35 with a new magnetic spectrometer led to the following values:

Card 2/3

Measurement of the Energy of the α-Particles of 8/056/60/039/01/11/029 an Emitter 8/056/8070

 $E_T = 4.394 \pm 0.002$  and  $E_{TT} = (4.213 \pm 0.002)$  MeV, and these agree very well with those obtained in the present work. There are 2 tables and 9 references: 1 Soviet, 1 South African, 2 Canadian, and 5 American.

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ASSOCIATION: Leningradskiy fiziko-tekhnicheskiy institut Akademii nauk SSSR (Leningrad Physicotechnical Institute of the Academy of Sciences, USSR)

SUBMITTED: March 22, 1960

Card 3/3

## 5/120/61/000/002/010/042

**AUTHORS:** 

Vorob'yev, A.A., and Korolev, V.A.

TITLE:

A method for measuring the transparency coefficient of

a grid in a pulse ionization chamber

PERIODICAL: Pribory i tekhnika eksperimenta, 1961, No. 2, pp. 78-80 Green's reciprocity theorem can be used to determine

the potential V+ induced on the collecting electrode by positive The latter is due to the fact that the grid introduced to screen the collector is not 100% efficient. The present authors point out that the theoretical treatment of the problem given by O. Buneman, T.E. Cranshaw and J.A. Harvey (Ref. 1; Canad. J. Res. A. 1949, 27, 191) is very complicated and, moreover, is based on various simplifying assumptions. They therefore suggest a method for the experimental determination of the grid transparency coefficient. Their argument runs as follows. The potential V+i induced by the i-th positive ion located at a distance x from the electrode 1 (Fig. 1) can be shown to be given by:

$$V^{+}_{,i} = \bullet \varphi (x)/CU$$
 (2)

Card 1/ 5

### S/120/61/000/002/010/042 E032/E114

A method for measuring the transparency coefficient of a grid in a pulse ionization chamber

where  $\phi(x)$  is the potential at the point x, which appears under the following conditions. The high-voltage electrode and the grid are earthed, the collecting electrode is at a potential U, and there is no space charge. The potential distribution in this case is shown in Fig.1. We are interested in the region between the high-voltage electrode and the grid. At a sufficient distance from the grid the field  $E_d$  is a constant, in which case

$$\varphi(\mathbf{x}) = \mathbf{E}_{\alpha}\mathbf{x} \tag{3}$$

The potential V+ is then given by

$$V^{+} = \sum_{i=1}^{N} V_{i}^{+} = Ne \widetilde{R} \cos \Theta E_{d}/CU$$
 (4)

where  $\overline{R}$  is the average distance of the ions from the beginning of the track and  $\theta$  is the angle between the direction of the track and the normal to the surface of the electrodes. In order to determine  $E_d$  the chamber is filled with pure argon and a high Card 2/5

S/120/61/000/002/010/042 E032/E114

A method for measuring the .... E032/E114

voltage is applied to the collecting electrode (U = 2 kV) while
the grid and the high-voltage electrode are earthed. The pressure
in the chamber is adjusted so that alpha particles emitted from a
source mounted on the high-voltage electrode stop within the
region in which the field is still constant. The field Ed which
penetrates beyond the grid extracts a fraction of electrons
produced by ionization and pulses appear on the collecting
electrode. Next, a compensating positive voltage AU is applied
to the high-voltage electrode and is adjusted until there are no
pulses at the collecting electrode. The pulses disappear when the
compensating field E' a and Ed are equal, i.e.

 $E'_{d} = E_{d} = \Delta U/d \tag{5}$ 

where d is the distance between the high-voltage electrode and the grid. Hence  $V^{+} = Ne \Delta U \tilde{R} \cos \theta/CU d \qquad (6)$ 

Thus, in order to determine  $V^+$  it is sufficient to carry out a simple experiment involving the determination of the compensating voltage  $\Delta U$ . The authors recommend pure argon as the working Card 3/5

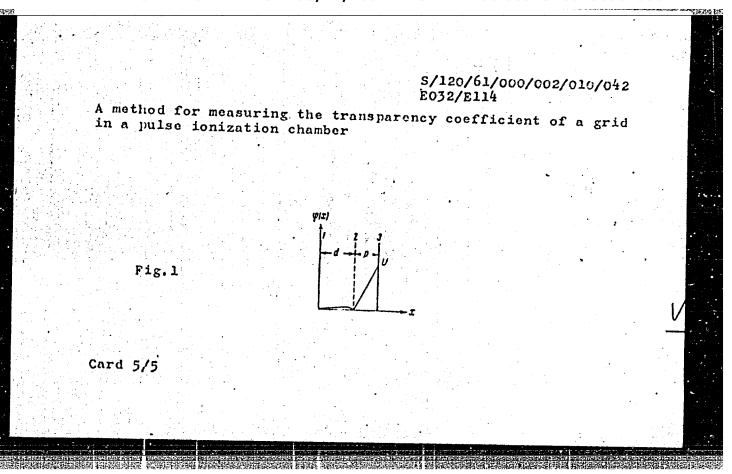
5/120/61/000/002/010/042

A method for measuring the transparency. E032/E114

gas. Best results are obtained by determining the compensation point using pulses from the high-voltage electrode, since in this case the sign of the pulses will change at the balance point. In this way an accuracy of 3-5% in  $\triangle$  U/U can be achieved. The experimental results are found to be in agreement with those computed by Buneman and Cranshaw and Harvey (Ref. 1) (see Table), the discrepancy between the experimental data and the theory being negligible in the case of grids with low transparency. Acknowledgements are expressed to A.P. Komar for interest in this work and to G.Ye. Solyakin for taking part in the discussions. There are 2 figures, 1 table and 2 references: 1 Soviet and 1 English.

ASSOCIATION: Leningradskiy fiziko-tekhnicheskiy institut (Leningrad Physico-technical Institute)

SUBMITTED:		March 24	1960		
grid	ı, mm	r, mm	r, mm	ocalc.	σ <sub>exp</sub> .
1	0.05	1.5	58	0.0065	0.0061
2 Card	0.05	3.0	58	0.0185	



s/120/61/000/004/003/~~4 E032/E514

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errementalism in the control of the

Vorob'yev, A.A. and Korolev, V.A.

A study of the properties of an argon-methane mixture AUTHORS: as the working gas of an ionization chamber TITLE:

PERIODICAL: Pribory i tekhnika eksperimenta, 1961, No.4, pp. 42-46

W. N. English and G. C. Hanna (Ref. 1: Canad. J. Phys., 1953, 31, 768) have recommended argon-methane mixtures as a suitable working gas for ionization chambers. The present authors report some measurements of the properties of such chambers. All the measurements were carried out with a plane ionization chamber containing a grid. The gas mixture was made up of commercial argon (Ar - 99.9%,  $O_2$  - 0.02%,  $N_2$  - 0.08%,  $CO_2$  - 0.005%) and commercial methane. The following quantities were measured: 1) Electron drift velocity, 2) the recombination rate, 3) the magnitude of the saturating field, 4) the electron attachment effect and the maximum permissible field, 5) the average energy of the electrons in the gas and 6) the dependence of the ionization on the energy of  $\alpha$ -particles traversing the chamber. It was found that with this gas mixture it is possible to choose the electrode potentials so that: Card 1/13

A study of the properties of ...

S/120/61/000/004/003/034 E032/E514

1) the recombination effects (between electrons produced during the ionization and the corresponding positive ions) are reduced to a minimum; 2) there is no appreciable attachment of electrons to neutral molecules and 3) the electrons can pass freely through the intermediate grid. The electron collection time can easily be reduced to 1-0.5 µsec and a plateau of several hundred volts is obtained, e.g. with a methane concentration of 5% the collection time is approximately 0.7 µsec. The ionization was found to be a linear function of  $\alpha$ -particle energy in the range 5.4-8.8 MeV. An important advantage of the argon-methane mixture is that it does not require any additional purification. Acknowledgments are expressed to A. P. Komar for his interest in this work and to M. F. Sobolevskaya for assistance with the measurements. There are 5 figures, 1 table and 8 references: 3 Soviet (1 a translation from English) and 5 non-Soviet. The following English-language references are given: Ref.1 (quoted in text), Ref. 5: C. E. Melton, G. S. Hurst, T. E. Bortner, Phys. Rev., 1954, 96, 643; Ref.6: G. Bertolini, M. Bettoni, A. Bisi, Phys. Rev., 1953, 92, 1586; Ref. 7: D. Strominger, T.M. Hollander, G. T. Seaborg, Rev. Mod. Phys., 1958, 30, 2. Card 2/4 3

A study of the properties of ... \$\frac{2959\lambda{l}}{5/120/61/000/004/003/034}\$

ASSOCIATION: Fiziko-tekimicheskiy institut AN SSSR (Physico-technical institute AS USSR)

SUBMITTED: March 24, 1960

Fig. 3: Amplitude V of palses at the collecting electrode on:
a - p.d. between high-voltage electrode and the grid and
6 - p.d. between the grid and the collector, (methane concentrations are marked on the curves.

8/020/61/136/004/008/026 B019/B056

AUTHORS:

Komar, A. P., Academician of the AS UkrSSR, Vorob'yev, A. A., and Korolev, V. A.

TITLE:

Measurement of the Fluctuation of Ionization Produced by α-Particles in Argon

PERIODICAL: Doklady Akademii nauk SSSR, 1961, Vol. 136, No. 4, PP · 795 - 797

TEXT: In the introduction, the authors refer to the frequently used measurement of ionization caused by nuclear particles for the purpose of determining the energy of nuclear particles. A relation given by V.Fano (Ref. 1) for the mean square fluctuation of the number of ion pairs with constant energy of the ionizing particles is written, and it is found that this formula is suited for determining the upper limit of the mean fluctuation, but not for more exact computations. Besides, Fano assumed that the ratio between the probabilities of the various inelastic processes is independent of the nature and energy of the ionizing particles. The measurements carried out by the authors were made by means Card 1/4

Measurement of the Fluctuation of Ionization S/020/61/136/004/008/026 Produced by  $\alpha$ -Particles in Argon S/020/61/136/004/008/026

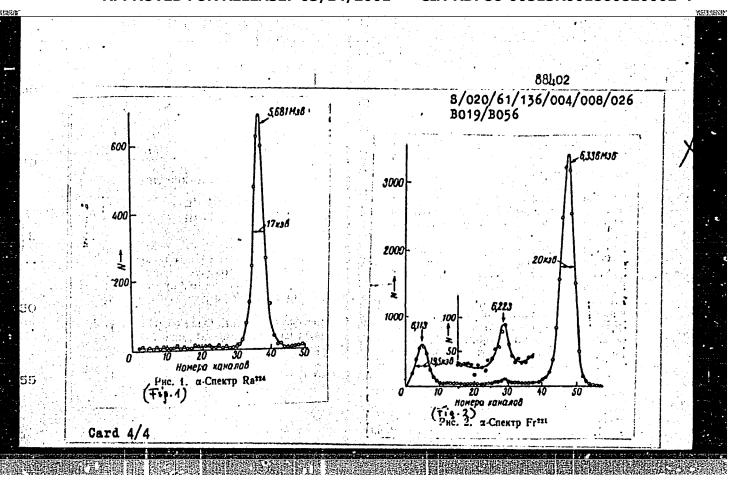
of  $\alpha$ -particles emitted by Ra <sup>224</sup> (E $_{\alpha}$  = 5.681 MeV) and of  $\alpha$ -particles emitted by Fr <sup>221</sup> (E $_{\alpha}$  = 6.336 MeV). The ionization chamber was filled with chemically pure argon + 1.5% CH $_{4}$ , whereby recombination could be prevented under certain conditions. Electronic collimation was used, whereby the resolution and, thus, the quality of the spectrum could be improved. The electronic means for improving the signal-to-noise ratio are briefly described. The measurements are graphically represented in Figs.1 and 2. The half-width of the Ra <sup>224</sup>  $\alpha$ -line is 17 keV and has a mean fluctuation of 7.2 keV. This mean fluctuation  $\delta$  is composed of  $\delta = \sqrt{\delta_{1}^{2} + \delta_{2}^{2} + \delta_{2}^{2}}$ , where  $\delta_{1}$ ,  $\delta_{2}$ ,  $\delta_{3}$  are the mean fluctuations which are due to the fluctuations of the ionization, to radio noise, and to other causes. In the case of Ra <sup>224</sup>,  $\delta_{3}$  is negligibly small, and because  $\delta_{1}$  = 4.7 keV, it follows that:  $\delta_{1}$  = 5.5 keV. For Fr <sup>221</sup>,  $\delta_{1}$  = 6.0 keV was obtained. From a discussion of the results, the authors conclude that  $\delta_{1}$  may be described by

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5	Measurement of the Fluctuation of Ionization S/020/61/136/004/008/026 Produced by α-Particles in Argon B019/B056	
10	$\delta_{\rm N}({\rm E}_{\alpha})=5.8/{\rm E}_{\alpha}/6.0$ (4) for different E <sub>\alpha</sub> . E <sub>\alpha</sub> must be given in Mev. In the relation $\delta_{\rm N}^2={\rm FN}_{\rm O}$ (1) given by Fano, where N <sub>O</sub> is the mean number of ion pairs, F is found equal to 0.22, and its upper limit is given as F <sub>lim</sub> = 0.33. The authors thank M. F. Sobolevskaya for her help in carrying out the measurements. There are 2 figures and 8 non-Soviet references: 5 US, 1 Canadian, 1 German, and 1 French.	
	ASSOCIATION: Fiziko-tekhnicheskiy institut Akademii nauk SSSR (Institute of Physics and Technology, Academy of Sciences USSR)	
	SUBMITTED: November 1, 1960	4
	Card 3/4	

"APPROVED FOR RELEASE: 03/14/2001 CIA-RDP86-00513R001860820002-4



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9,91	s/020/61/137/001/009/021 B104/B209	
26.2J	Vorob!yev. A. A., Komar, A. P., Academician AS UkrSSR, and Korolev, V. A.	in the
TITLE:	The possibilities of reducing the effect of ionization fluctuations in gases	
PERIODICAL:	Doklady Akademii nauk SSSR, v. 137, no. 1, 1961, 54-57	15
Phys. Rev., mean square energy of th	uthors based their work on a paper by Fano (Ref. 1: U. Fano, 72, 26 (1947)), in which an expression was obtained for the fluctuations of the number of ion pairs at a constant e ionizing particles. Fano's calculations show that these are determined chiefly by the redistribution of ionized are determined their total amount fluctuates less.	20
and excited The authors	atoms. Evidently, their total have now determined the amount of fluctuations of the total have now determined the amount of fluctuations of the total	25
ed the mean	square fluctuation $\delta_J^2$ of the total ionization $\overline{J}$ :	
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S/020/61/137/001/009/021 B104/B209

The possibilities of reducing ...

 $\delta_J = \frac{A^2}{N_0 P} \left(n_k - \frac{E_k}{W}\right)^2 = \frac{F}{N_0} . \quad N_0 \text{ denotes the mean number of ion pairs,}$   $W = W_0/(1 + \sigma(1-P)/P) = W_0 A, \quad P = \sum P_k^i \text{ the total probability of ionization in inelastic collision, } W_0 \text{ the mean energy of ion pair production without additional ionization, and } n_k \text{ the number of ions}$  produced in the k-th collision. The relations

$$F = \Phi(s) + \frac{1}{PW_0^2} \Big[ \sum_{k} P_k^l (W_l - E_k^l)^2 + \sum_{k \in SG} P_k^l (W_l - E_k^l)^2 \Big]; \tag{8a}$$

$$\Phi(\sigma) = \frac{1}{W_0^2} \left[ (W - W_i)^2 + \sigma \frac{1 - P}{P} (W - W_e)^2 + \frac{1 - P}{P} (1 - \sigma) W_e^2 \right]. \quad (86)$$

are obtained for F. The last two terms in (8a) are due to fluctuations of the energy losses during ionization and excitation, and do not depend

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The possibilities of reducing ...

on the additional ionization.  $\Phi(\sigma)$  is determined by the redistribution of the number of ionized and excited atoms, as well as by the fluctuations of the number of ionized and excited atoms, as well as by the fluctuations arising in the additional ionization. In the limiting case where arising in the additional ionization is missing  $(\sigma = 0)$ , Eq. (8a) goes over into the additional ionization is missing  $(\sigma = 0)$ , Eq. (8a) goes over into the formula of Fano. Fig. 1 shows the ratio  $\Phi(\sigma)$  as depending on the

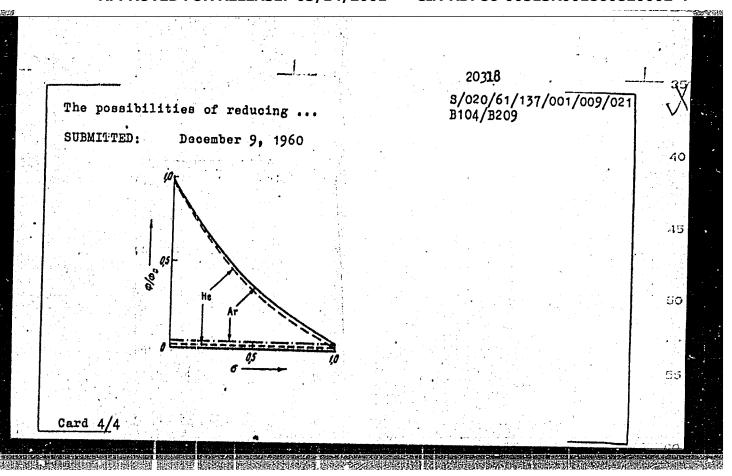
probability  $\sigma$  of additional ionization for He and Ar. It is seen that  $\Phi(\sigma)$  for argon drops to nearly one-thirtieth with rising probability, and for helium it drops to nearly one-hundredth. The first of the terms appearing in (8a) was found to be always about 0.03, and the second is negligible. From this it follows that the accuracy of measurement of the energy of ionizing particles is considerably improved by recording all ionized and excited atoms. There are 1 figure and 3 non-Soviet-bloc references.

ASSOCIATION:

Fiziko-tekhnicheskiy institut Akademii nauk SSSR (Institute of Physics and Technology of the Academy of

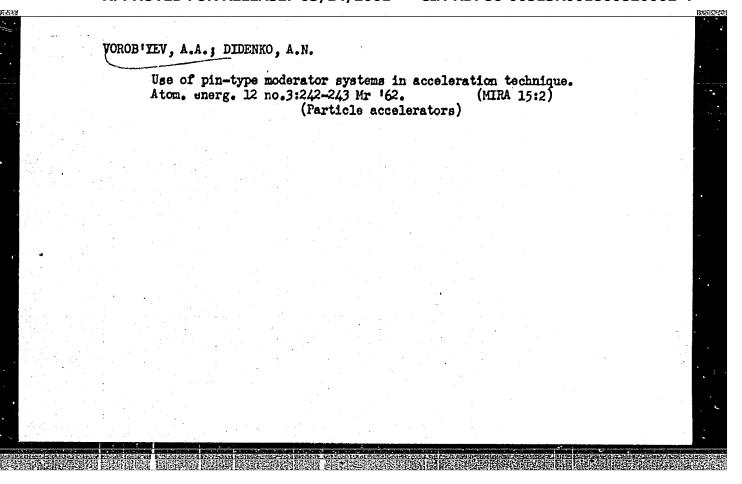
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### "APPROVED FOR RELEASE: 03/14/2001 CIA-RDP86-00513R001860820002-4



39479 5/056/62/043/002/005/053

Vorob'yev, A. A., Komar, A. P., Korolev, V. A. 26 2311

AUTHORS: Decrease of ionization fluctuations of  $\alpha$ -particles in argon

TIPLE: Zhurnal eksperimental noy i teoreticheskoy fiziki, v. 43,

PERIODICAL: no. 2(8), 1962, 426-428

TEXT: The authors had shown earlier (DAN SSSR, 137, 54, 1961) that the ionization fluctuations associated with redistributions of the numbers of excited and ionized molecules can be reduced by adding a gaseous impurity with an iorization potential lower than the energy of the lowest excited level of the principal component. Here, the authors tried to check this possibility by experiment. They used a pulsed ionization chamber filled with argon containing 0.17 % N<sub>2</sub>, 0.02 % O<sub>2</sub>, and an acetylene impurity. As

its ionization potential of 11.35 ev is lower than the lowest brgon level (11.5 ev), the acetylene addition increases the ionization. The ionization fluctuations were calculated from the half-width of the  $\alpha$ -line

 $(E_{\alpha} = 5.681 \text{ MeV})$  of Ra<sup>224</sup>; for comparison, the measurements were repeated

Card 1/2

Decrease of ionization ...

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on Ar + 1 % CH2.

 $\Delta E_{\alpha}$ , kev  $\Delta E_{fl}$ , kev  $\Delta E_{N}$ , kev 8.1 5.8 5.7 216 215000 0.22 Ar+0.8% C2H2 6.0 4.7 3.7 175 268000 0.05

where  $\Delta E_{\alpha}$  = total root-mean-square pulse-height fluctuations,  $\Delta E_{T_{1}}$  = rootmean-square pulse-height fluctuations due to electronic noise,  $\Delta E_{\rm N}$  = the same due to fluctuations in the number of ion pairs, N = total number of ion pairs,  $\Delta N = \text{root-mean-square}$  fluctuation in the number of ion pairs; F is determined by  $\Delta N/N = \sqrt{F/N_0}$ ;  $N_0 = 212,000$  ion pairs. The maximum halfwidth of the  $\alpha$ -line was 8.7 kev. There are 1 figure and 1 table.

ASSOCIATION: Fiziko-tekhnicheskiy institut im. A. F. loffe Akademii nauk SSSR (Physicotechnical Institute imeni A. F. Ioffe of the Academy of Sciences USSR)

SUBMITTED:

March 13, 1962

Card 2/2

8/0139/64/000/001/0173/0177

ACCESSION NR: AP4020312

AUTHOR: Vorob'yev, A. A.

TITLE: Investigations in betatron physics reported in the 1962 Bucharest colloquium on construction and application of betatrons

SOURCE: IVUZ. Fizika, no. 1, 1964, 173-177

TOPIC TAGS: betatron, bremestrahlung radiation, accelerated electron, radiation intensity, equilibrium orbit, field decay, air gap, electromagnet, beam focusing, betatron TPI

ABSTRACT: The highlights of the colloquium on betatron physics held in Bucharest during 14-20 November 1962 are reported. Soviet Russia, Germany, Yugoslavia, Poland, Rumania, and Hungary participated. Some noteworthy papers were presented by Professor E. Eckard of Jena University, Professor A. A. Vorob'yev of Tomsk Polytechnic Institute, and Engineer K. Iliyesku. Professor Ekkard spoke on improved betatron operation reliability, improved bremsstrahlung radiation from targets, regulation, measurement and stability of betatron end energy, and accelerated electron input beams. A detailed description is given of the 25 Mev Rumanian betatron with 42 roentgen/m/min radiation intensity, equilibrium orbit radius of 250 mm,

Card 1/2

# field decay index of 0.75, air gap of 67.4 x 78 mm, electromagnet weight of 3 tons, supply voltage of 5.3 kv, and vacuum of 1 to 5 torr. The ceramic chamber is shown to have an elliptic cross section. Several papers by Iliyesku and his co-workers

to have an elliptic cross section. Several papers by Iliyesku and his co-workers included topics on magnetic field characteristics, weak and strong beam focusing techniques, air gap studies, accurate field decay measurements, and improved betatron emission. Geske and Schmaltz of the German Democratic Republic outlined the design of a 30 Nev betatron under construction. A. Brinshek of Yugoslavia reported on the automatic matching technique for supply voltage and frequency in the Yugoslavian 31-Mev betatron. Burger, Rudolph, Plitz, Koshma, Lungu, Simionescu, I. Leybovich, S. Haltrich, V. Grechescu, Klapper, Bezich, Yamnik, Kernel, Schmeider, Vaksel', Shtern also participated. Various problems for future consideration were discussed. Orig. art. has: 2 figures.

ASSOCIATION: Tomskiy politekhnicheskiy institut imeni S. M. Kirova (Tomsk Polytechnic Institute)

SUBMITTED: 06May 63

ACCESSION NR: AP4020312

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Card 2/2

### "APPROVED FOR RELEASE: 03/14/2001

CIA-RDP86-00513R001860820002-4

ACCESSION NR: AP4036575

AUTHOR: Vorob'yev, A. A.

TITLE: Solective mechanical failure under acoustic compression or dilation waves on composition of materials having different hardnesses

SOURCE: IVUZ. Fizika, no. 2, 1964, 168-169

TOPIC TAGS: dynamic load, modulus of elasticity, mechanical failure, acoustic wave, compression wave

ABSTRACT: The author discusses the fact that under dynamic loads the failure rate in stronger substances is greater than in weaker substances. Consequently, the clastic wave propagation (compression or dilation) is faster in the stronger substances (with a higher modulus of elasticity) than in the weaker ones. It is argued that when a material of compound composition, comprising a mixture of strong and weak zones, is subjected to dynamic loading, its weak parts will not reach the failure stage before the load on the stronger zones exceeds the elastic limit and causes complete failure of these stronger zones. It is believed that

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ACCESSION NR: AP4033111

5/0120/64/000/002/0069/0071

AUTHOR: Alkhazov, G. D.; Yorob'yev, A. A.; Korolev, V. A.; Seliverstov, D. M.

TITLE: Simple counting unit for a slow multichannel analyzer

SOURCE: Pribory\* i tekhnika eksperimenta, no. 2, 1964, 69-71

TOPIC TAGS: pulse height analyzer, multichannel pulse height analyzer, slow multichannel analyzer, pulse counter, pulse counting unit

ABSTRACT: A simple counting unit intended for low (up to 100 pulse/sec) counting rates is described. The unit is recommended for long-duration few-pulse applications where reliability is the main requirement. The 32-tube counting unit was used in a 20-channel time analyzer and tested for several months under actual operating conditions. The a-c line supply voltage is stabilized by a ferroresonance stabilizer only; variation of the plate voltage

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ACCESSION NR: AP40331			
within 200-250 v does not z	esult in false operation. The ction diagram is shown and ex	power consumption is iplained in the article.	
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5/0141/64/007/002/0338/0342

ACCESSION HR: AP4039734

AUTHOR: Vorob'yev, A. A.; Bezmaterny\*kh, L. N.; Didenko, A. N.; Lisitsy\*n, A. I.; Ol'shanskiy, A. P.

TITLE: Laminated dielectric coatings with large reflection coefficients

SOURCE: IVUZ. Radiofizika, v. 7, no. 2, 1964, 338-342

TOPIC TAGS: dielectric coating, reflection coefficient, cavity resonator, microwave equipment, dielectric permittivity

ABSTRACT: In view of various applications of laminated dielectric coating with large reflection coefficients, their reflecting properties are analyzed on the basis of a calculation of the reflection coefficient from a semi-infinite periodic medium, comprising an infinite waveguide of arbitrary cross section, one half of which is filled with dielectric layers. Such a representation neglects the reflection from the second boundary of the layer and is justified at the frequencies considered. The field outside the outermost layer is then described as a sum of incident and reflected waves, and inside the layer by a wave traveling inside the dielectrics. Calculations show that for a given reflection coefficient the dimensions of the laminated coating decrease sharply with increasing dielectric con-

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### ACCESSION NR: AP4039734

stant of the layers, and in the case of large dielectric constants (e. g., barium titanate), such layers can be used not only in the optical but also in the microwave bands. It is shown that a frequency exists at which the tangential electric field on the surface of the laminated medium vanishes, making it possible in some cases to replace metallic walls of cavity resonators by laminated dielectrics without distorting the field structure in the cavity. Tests, of laminated dielectric consisting of alternating layers of paraffin and foamed plastic placed in a rectangular waveguide confirmed this assumption, and the cavity produced by shorting the ends of this waveguide had approximately the same Q as a metal cavity. Slight deviations from theory are explained. The use of dielectrics with large permittivities  $(10^2-10^3)$  will make it possible to reduce the total thickness of the sandwich to 1-2 cm in the 10-cm band and to several tenths of a millimeter in the millimeter band. Orig. art. has: 2 figures and 11 formulas.

ASSOCIATION: None

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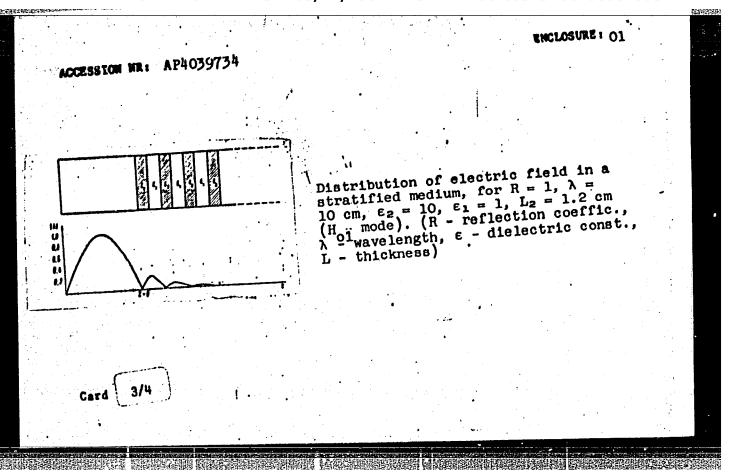
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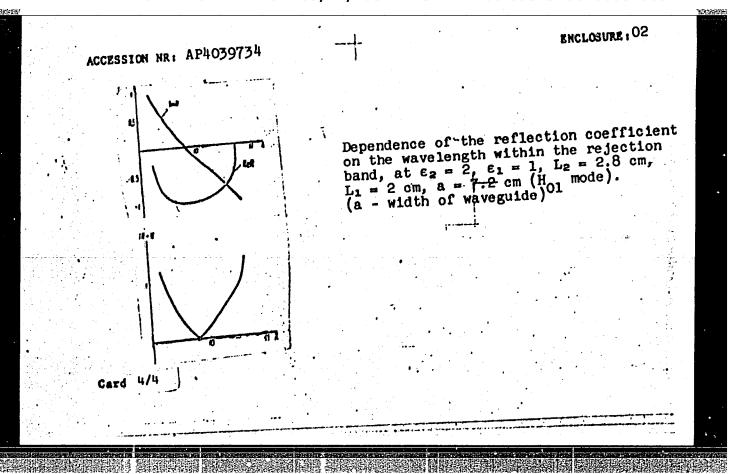
SUB CODE: EM, MT

NR REF SOV: 003

OTHER: 04

Card 2/4





- ACC NR. AP7005707

SOURCE CODE: UR/0089/67/022/001/0003/0006

AUTHOR: Vorob'vev. A. A.; Didenko, A. N.; Ishkov, A. P.; Kolomenskiy, A. A.; Lebedev, A. N.; Yushkov, Yu. G.

ORG: none

TITIE: Investigation of autoresonant method of particle acceleration by electromagnetic waves

SOURCE: Atomnaya energiya, v. 22, no. 1, 1967, 3-6

TOPIC TAGS: particle acceleration, magnetic resonance, electron waveguide, electron accelerator, ELECTROMAGNETIC CUAVE

ABSTRACT: This is a continuation of earlier work (in: Trudy Mezhdunarodnoy konferentsii po uskoritelyam [Trans. Internat. Conf. on Accelerators] (Dubna, 1963), M., Atomizdat, 1964, p. 1030, and earlier papers) which demonstrated the feasibility of resonant acceleration of particles by a transverse wave in a longitudinal magnetic field under suitable conditions. The present paper contains the results of an experimental investigation of this method of acceleration. In view of the limited possibility of obtaining the required strong field in a large volume, the study is confined to acceleration by 10-cm electromagnetic waves inside a straight smooth waveguide (H11 and H10 modes). Equations are derived for the angular velocity and phase of a particle accelerated in such a structure, and for the length of the waveguide over which the particle energy will increase. The accelerating system was a

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UDC: 621.384.62

' ACC NR: AP7005707 rectangular waveguide (72 x 44 mm) for the H10 mode or a 76-mm diameter round waveguide for the H11 mode. The length of the waveguide ranged from 150 to 1000 mm. Pulsed microwave power (not more than 600 kw) (3000 MHz) was fed in 3-usec pulses at a repetition frequency of 50 Hz. The 1000-oe dc field was produced with a solenoid. The particle energy was determined from the deceleration produced by aluminum folls and reached 700 kev, at an electric field intensity of 3-5 kv/cm, which is higher than obtainable by ordinary cyclotron acceleration. The ancillary tests made on the equipment are briefly described. The experimental data agree with the earlier theoretical predictions and it is concluded that the autoresonant mechanism can be used for effective injection of particles into magnetic traps. Orig. art. has: 1 figure and 9 formulas. SUB CODE: 20/ SUBM DATE: 05Sep66/ ORIG REF: 006/ ATD PRESS: 5117 Card

ACC NR: AT7003986

SOURCE CODE: UR/0000/66/000/000/0005/0010

AUTHOR: Vorobiyev, A. A.; Kalganov, A. F.; Lukutin, V. A.; Patsevich, V. V.

ORG: Tomsk Polytechnic Institute (Tomskiy politekhnicheskiy institut)

TITLE: Theory and technology of electrostatic machines

SOURCE: Mezhvuzovskaya konferentsiya po elektronnym uskoritelyam. 5th, Tomsk, 1964. Elektronnyye uskoriteli (Electron accelerators); trudy konferentsii. Moscow, Atomizdat, 1966, 5-10

TOPIC TAGS: electrostatic generator, particle acceleration, electronic test

ABSTRACT: The phenomena transpiring in the electrostatic generator and their analogy to the phenomena in the electromagnetic generator are briefly reviewed (e.g., D. Gignoux, "Electrostatic generators for space application", 102-ème Colloque du SNRC, Grenoble, 1960). Formulas for maximum power of disk-type and cascaded-conveyer generators show that the maximum specific power (per unit volume or weight) is inversely proportional to the stator-rotor gap; the load voltage and current are independent of the gap. Small gaps are preferable because they mean smaller spurious capacitance, and the available power becomes closer to its theoretical value. The latter statement was proved theoretically and experimentally, on a single-disk generator, at the NII of Nuclear Physics, Tomsk Polytechnic Institute. An electrostatic generator with parallel-connected poles and vacuum insulation seems to be most promising. Orig. art. has: 8 formulas.

Cord 1/1 SUB CODE: 09 / SUBM DATE: 06Mar66 / ORIO REF: 003 / OTH REF: 003

ACC NR: AT7003994

SOURCE CODE: UR/0000/66/000/000/0075/0082

AUTHOR: Vorob'yev, A. A.; Bezmaternykh, L. N.; Didenko, A. N.; Filatova, R. M.

ORG: Scientific Research Institute of Nuclear Physics, Electronics, and Automation, Tomsk Polytechnic Institute (Nauchno-issledovatel'skiy institut yadernoy fiziki, elektroniki i avtomatiki pri TPI)

TITLE: Waveguide accelerating systems with walls not shielding the control magnetic field

SOURCE: Mezhvuzovskaya konferentsiya po elektronnym uskoritelyam. 5th, Tomsk, 1964. Elektronnyye uskoriteli (Electron accelerators); trudy konferentsii. Moscow, Atomizdat, 1966, 75-82

TOPIC TAGS: waveguide, accelerator, cyclic accelerator, particle acceleration

ABSTRACT: A multilayer-dielectric coating similar to that used in Fabry-Perot interferometers (W. Gulshaw, Proc. Phys. Soc., London, v. 66, sec. B, 597, 1953) and in lasers (J. Franklin Inst., 273, 177, 1962) is proposed for the walls of waveguide-type accelerators. Uniformly bent smooth and septate closed

Card 1/2

ACC NR: AT7003994

rectangular waveguides with multilayer-dielectric walls are theoretically and experimentally investigated. Formulas for the rejection frequency of a periodic multilayer structure, for attenuation, and for the total electromagnetic-wave losses due to reflection from a multilayer dielectric are derived. A length of standard  $72 \times 34$ -mm waveguide whose ends were closed by multilayer-dielectric walls was excited by TE<sub>101</sub> -mode at  $\lambda = 10$ , 182 cm; at room temperature, Q = 1800. Findings: (1) At a fixed frequency, the field structure in the above system does not differ from that in an all-metal system; (2) Use of TE-modes is preferable; inside the multilayer wall, the field attenuates rapidly; with proper selection of wall parameters, no hazard of dielectric breakdown by SHF high power will'exist; (3) The above multilayer-dielectric walls are feasible if Sr and Ba titanates are used as materials (see R. O. Bell et al., IRE Trans., MTT-9, 239, 1961). Orig. art. has: 3 figures, 15 formulas, and 1 table.

SUB CODE: 09 / SUBM DATE: 06Mar66 / ORIG REF: 001 / OTH REF: 003

Card 2/2

ACC NR: AT7003986

SOURCE CODE: UR/0000/66/000/000/0005/0010

AUTHOR: Vorob'yev, A. A.; Kalganov, A. F.; Lukutin, V. A.; Patsevich, V. V.

ORG: Tomsk Polytechnic Institute (Tomskiy politekhnicheskiy institut)

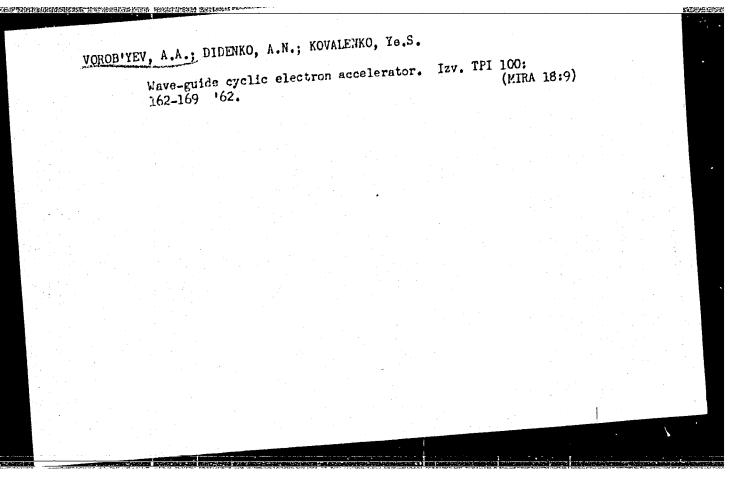
TITLE: Theory and technology of electrostatic machines

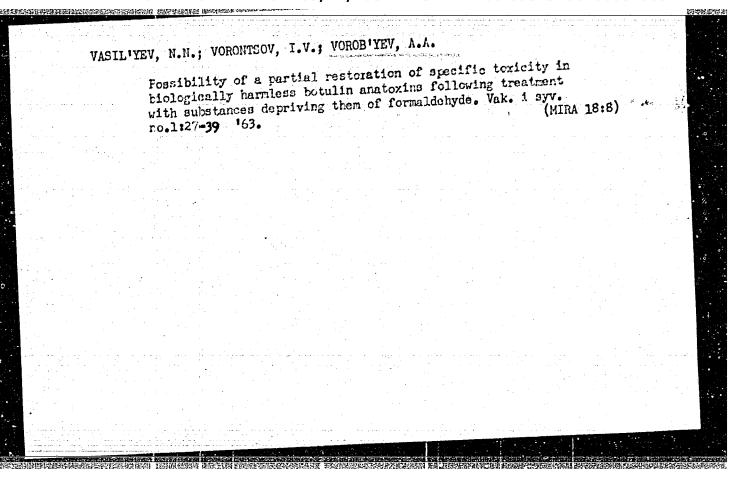
SOURCE: Mezhvuzovskaya konferentsiya po elektronnym uskoritelyam. 5th, Tomsk, 1964. Elektronnyye uskoriteli (Electron accelerators); trudy konferentsii. Moscow, Atomizdat. 1966. j-10

TOPIC TAGS: electrostatic generator, particle acceleration, electronic test

ABSTRACT: The phenomena transpiring in the electrostatic generator and their analogy to the phenomena in the electromagnetic generator are briefly reviewed (e.g., D. Gignoux, "Electrostatic generators for space application", 102-ème Colloque du SNRC, Grenoble, 1960). Formulas for maximum power of disk-type and cascaded-conveyer generators show that the maximum specific power (per unit volume or weight) is inversely proportional to the stator-rotor gap; the load voltage and current are independent of the gap. Small gaps are preferable because they mean smaller spurious capacitance, and the available power becomes closer to its theoretical value. The latter statement was proved theoretically and experimentally, on a single-disk generator, at the NII of Nuclear Physics, Tomsk Polytechnic Institute. An electrostatic generator with parallel-connected poles and vacuum insulation seems to be most promising. Orig. art. has: 8 formulas.

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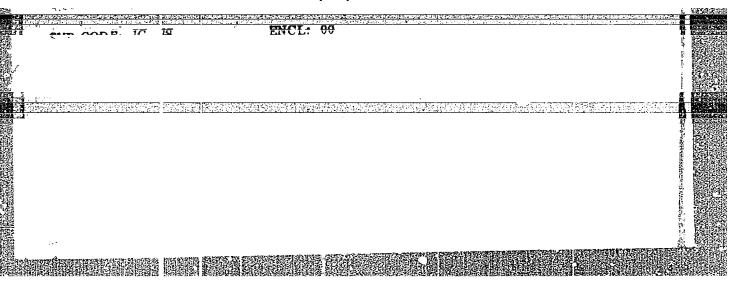
COROBITEV, A.A.; DIDENKO, A.N.; LISITSYN, A.I.; MOROZOV, B.N.; POTEKHIN, Yu.I.; SALIVON, L.G.; FILATOVA, R.M.

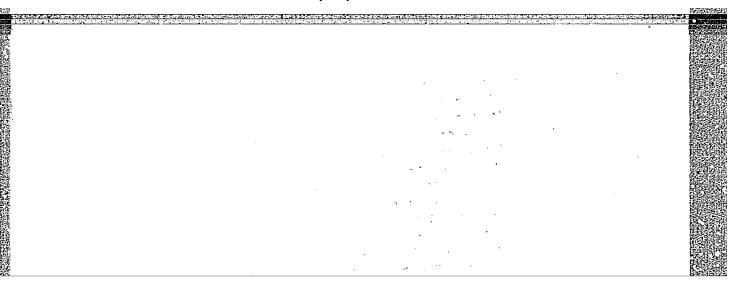
A 10-Mev. wave-guide type synchrotron. Atom. energ. 18 no.6:633-634
Je '65.

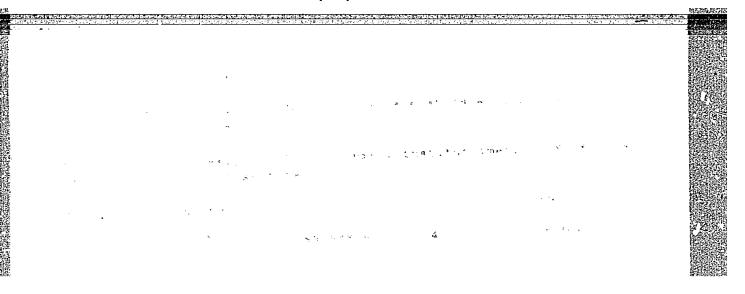
(MIRA 18:7)

VOROB!YEV, Anatoliy Andreyevich; VASIL!YEV, Nikolay Nikolayevich;
KRAVCHENKO, Anatoliy Timofeyevich; ANAN!YEV, V.A., red.

[Anatoxins] Anatoksiny. Moskva, Moditsina, 1965. /67 p.
(MIRA 18:10)







VOROB'YEV, A.A.; VOROB'YEV, G.A.; KOCHERBAYEV, T.K.; KOSTRYCIN, V.A.; NEKRASOVA,

Effect of electrodes and the structure of a dielectric crystal on its electric strength. Fiz. tver. tela 6 no.5:1560-1562
My '64. (MIRA 17:9)

1. Tomskiy politekhnicheskiy institut imeni Kirova.

VOROBLYEV, A. A., doktor fiziko-matematicheskikh nauk, prof.; VOROBLYEV, G. A., kand. tekhn. nauk; KOSTRYGIN, V. A., kand. tekhn. nauk

Dependence of the electrical strength of solid dielectrics on the thickness of the breakdown layer. Izv. vys. ucheb. zav.; energ. 7 no.5:108-110 My 164. (MIRA 17:7)

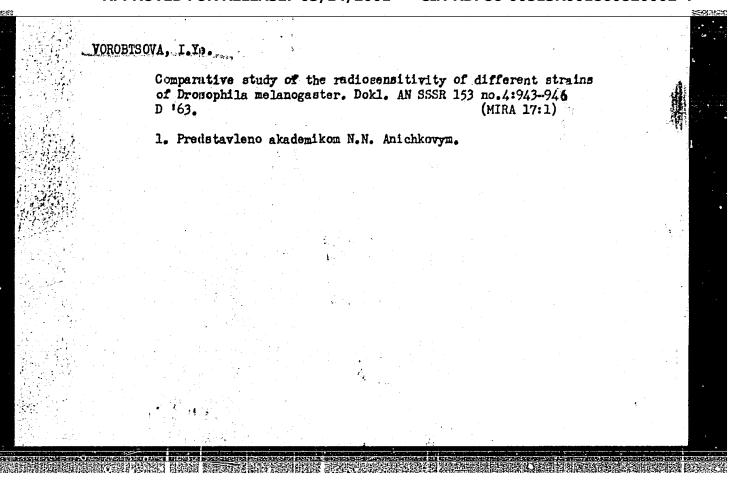
1. Tomskiy ordena Trudovogo Krasnogo Znameni politekhnicheskiy institut imeni Kirova. Predstavlena kafedroy tekhniki vysokikh napryazheniy.

L 36481-66 EWT(1) ACC NR: AP6027071 SOURCE CODE: UR/0139/66/000/002/0153/0155 AUTHOR: Vorobyev. A. A. 62 ORG: Tomsk Polytechnical Institute im. S. M. Kirov (Tomskiy politekhnicheskiy \_3 institut) TITIE: Remarks on the free energy and the lattice energy of real ionic crystals SOURCE: IVUZ. Fizika, no. 2, 1966, 153-155 TOPIC TAGS: ionic crystal, free energy, crystal lattice energy, solid physical property, quantum theory, thermodynamics ABSTRACT: The author considers the linear relation between the free energy and the lattice energy of real ionic crystals and experimentally establishes the relations between the various properties of such crystals. Normal mechanical, thermal, optical, and electrical properties of solids are described by classical thermodynamic and quantum theory with the aid of the free energy. It is shown that the lattice energy can be used instead of the free energy, and results are in good agreement with the conclusions of the dynamic theory of crystal lattices. The relations between the various properties of ionic crystals are then established in terms of the lattice energy, and it is pointed out that this affords a means of comparing the variations in the properties of ionic compounds due to variations in composition or experimental conditions. Lattice energy can thus serve as a single parameter to classify various ionic compounds and their mechanical and thermal properties. The author thanks Docents A. S. Naumova and V. A. Zhdanov for reviewing the paper and for useful remarks. Orig. art. has: 9 formulas. [JPRS: 36,364] SUB CODE: 20 / SUBM DATE: 16Nov64 / ORIG REF: 006 Card 1/1 mlf 

EWP(e)/EWT(m)/EWP(v)/T/EWP(t)/ETI/EWP(k)IJP(c) JD/HM/HW AT6024936 (A,N) ACC NR SOURCE CODE: UR/2981/66/000/004/0238/0253 AUTHOR: Kovrizhnykh, V. G.; Vorob'yev, A. A.; Ponogaybo, Yu. N.; Tsabrov, N. D.; Matveyev, B. I. B+1 ORG: none TITIE: Preparation of weldable sheets of SAP-1 alloy by coil rolling SOURCE: Alyuminiyevyye splavy, no. 4, 1966. Zharoprochnyye 1 vysokoprochnyye splavy (Heat resistant and high-strength alloys), 238-253 TOPIC TAGS: sintered aluminum powder, hot rolling, cold rolling, sheet metal ABSTRACT: The purpose of the work was to determine the feasibility of preparing thin sheets 0.6 to 3 mm thick of industrial dimensions (1000-1400 mm wide and 3500-7000 m long) from fusion-welded SAP-1 material (a sintered aluminum powder material) by coil rolling on existing industrial equipment, and also to study the mechanical properties and structure of hot- and cold-rolled sheets in relation to the conditions of deformation tion and annealing. It was found possible to produce such sheets by using a billet made by stamping on a vertical hydraulic press, and to weld them by fusion. Vacuum annealing can be replaced by long high-temperature annealing without vacuum for the purpose of adequately degassing the briquet and imparting weldable properties to the SAP-1 material. In order to obtain the maximum strength characteristics at high temperatures, the sheets should be produced only by hot rolling, If thin sheets cannot Card

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be produced by	hot rolling	alone, the cold rolling should be Orig. art. has: 6 figures and 3 t	carried out with	a mini-
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VOROB'YEV; MAKSIMOV; BEZRUKIKH, P.

Replies to the editors. Sots. trud 7 no.9:145-146 S '62.

(MIRA 15:9)

1. Zamestitel' predsedatelya pravleniya Gosudarstvennogo banka
SSSR (for Vorob'yev). 2. Zamestitel' nachal'nika otdela truda
SSSR (for Waksimov). 3. Nachal'nik
i zarabotnoy platy Gosplana SSSR (for Maksimov). 3. Nachal'nik
upravleniya bukhgalterakogo ucheta i otchetnosti Ministerstva
finansov SSSR (for Bezrukikh).

(Wages)

8/0189/63/000/003/0048/0051

AUTHORS: Vorobiyev. A. F.; Monayenkova, A. S.; Privalova, N. M.; Skuratov, S. M. ACCESSION NR: TITLE: Enthalpy of formation for OH, Kt, Hat, Lit ions in water solutions SOURCE: Moscow. Universitet. Vestnik. Seriya 2, Khimiya, no. 3, 1963, 48-51 TOPIC TAGS: alkali, enthalpy of formation, hydroxyl ion, cation, dilute solution

ABSTRACT: The enthalpy of formation for OH and alkaline ions has been calculated at 25C using international tables, 1961 edition, for the atomic weights. A root mean square error analysis has been made to determine the limits of calculation accuracy. For the hydroxyl ion, the calculation is based on the neutralization enthalpy of a strong acid by a strong base in an infinitely dilute solutions OH (p-pos) + H(p-pos) = H2O(liquid) . This leads to the values

 $\Delta H_{r} (OH_{ppm}) = \Delta H_{r} (H_{s}O_{11q}) - \Delta H_{neut} \Delta H_{r} (H_{ppm}) = (-68.32 \pm 0.01) - (-13.34 \pm 0.02) - O = -54.98 \pm 0.02 \text{ kcal/g-ion}$ . When the enthalpy of formation of the hydroxyl ion is known, the A Hp of the cation alkali metals can be determined from the known Alp of alkali-hydroxides in an infinitely dilute

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olution.	In an equat	ion form th	is yields AH.	$P^{(M_{p-p oo})} = 0$	Hy (MOHppe)	22-24-2	
= Na. Li. s listed	K. J. Final	ly, the ent	halpy of form	nation for w	H and the th Li+ 0,02 -66,58±0.0		<b>5</b>
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